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Excitation strengths in $^{109}\text{Sn}$: Single-neutron and collective excitations near $^{100}\text{Sn}$

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A set of $B(E2)$ values for the low-lying excited states in the radioactive isotope $^{109}\text{Sn}$ were deduced from a Coulomb excitation experiment. The 2.87-MeV/u radioactive beam was produced at the REX-ISOLDE facility at CERN and was incident on a secondary $^{58}\text{Ni}$ target. The $B(E2)$ values were determined using the known $2^+ \rightarrow 0^+$ reduced transition probability in $^{58}\text{Ni}$ as normalization with the semiclassical Coulomb excitation code GOSIA2. The transition probabilities are compared to shell-model calculations based on a realistic nucleon-nucleon interaction and the predictions of a simple core-excitation model. This measurement represents the first determination of multiple $B(E2)$ values in a light Sn nucleus using the Coulomb excitation technique with low-energy radioactive beams. The results provide constraints for the single-neutron states relative to $^{100}\text{Sn}$ and also indicate the importance of both single-neutron and collective excitations in the light Sn isotopes.

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We report here on measurements using Coulomb excitation of the nucleus $^{109}\text{Sn}$, with the aim of investigating the nature of the low-lying excited states in this nucleus. Of particular interest is the location of the single-neutron states, which provide important constraints for shell-model calculations. For the light Sn isotopes, these include the $g_{7/2}$, $d_{5/2}$, $s_{1/2}$, $d_{3/2}$, and $h_{11/2}$ single-neutron orbits. Only a limited number of studies have been reported using Coulomb excitation in the odd-Sn isotopes. In earlier work, the technique was used to investigate the collectivity of excited states in $^{115,117,119}\text{Sn}$ [1,2]. More recently, we reported on the first Coulomb excitation measurement in the isotope $^{107}\text{Sn}$ [3]. That measurement resulted in an enhanced reduced transition probability in comparison to shell-model calculations, which indicates the importance of mixing between single-neutron and collective excitations near the doubly magic nucleus $^{100}\text{Sn}$. In addition, in a series of recent measurements, Coulomb excitation was used to determine the reduced transition probabilities of the first $2^+$ states in the neutron-deficient Sn isotopes $^{106,108,110}\text{Sn}$ [4–7]. Also there the reduced transition probabilities were increased by ~30% with respect to state-of-the-art shell-model calculations based on realistic nucleon-nucleon interactions.

The experimentally known systematics for the single-neutron states in the odd-Sn nuclei are highlighted in Fig. 1 [8]. In $^{111}\text{Sn}$, these states have been firmly identified in previous reaction studies [9]. However, similar data are not available for the lighter nuclei and the current information is based on a number of different measurements, comparison to theory, and systematic trends. Down to $^{103}\text{Sn}$, the ground-state and first excited-state spins are believed to be $5/2^+$ and $7/2^+$, respectively. Recently, the energy splitting between these two levels in $^{105}\text{Sn}$ has been reported in two different measurements as 172 keV [10,11]. The energy splitting is similar to the 168-keV difference in $^{103}\text{Sn}$ [12]. Based on the energy systematics alone, this favors a $5/2^+$ ground state and a
7/2+ first excited state. However, measurements in Ref. [11] indicated that the level ordering is instead inverted with respect to 103Sn. This could be the result of a strong pairing interaction indicated that the level ordering is instead inverted with respect to the neutron-deficient Sn nuclei, using energies given in Ref. [8].

In addition to states of single-neutron character, we might also expect to observe states originating from the coupling of a particle or hole to a core state in a neighboring even-even Sn isotope [17]. These excitations are expected to have larger electromagnetic properties that were good examples of single-neutron and collective excitations [1]. Here, we aim to extend this type of investigation to the nucleus 109Sn.

The 109Sn radioactive beam was produced at the REX-ISOLDE facility at CERN. The method has been described in detail in a number of previous publications [3,5,7] and is presented only briefly here. The isobaric beam was produced by reactions in a 27-g/cm² LaC₃ target induced by a 1.4-GeV proton beam with a maximum intensity of ∼3 × 10¹³ protons per pulse (typically 9 pulses/min), ionization using a resonant laser ionization scheme [18], and separation according to A = 109 in the general purpose separator. The radioactive beam, consisting of Sn and In, was accelerated to an energy of 2.87 MeV/n before bombarding the 99.93% isotopically enriched 58Ni secondary target. The beam purity was calculated to be 89(1)% from the number of particles collected in a double-sided silicon strip detector (DSSSD) placed downstream of the secondary target [19] from the laser-on and laser-off runs. The Sn beam was estimated to have an intensity of ∼10⁹ particles per second.

Following Coulomb excitation in the secondary target, the emitted γ rays were identified using the MINIBALL γ-ray detector array [20]. The Ge clusters of the array were positioned around the target. The DSSSD was used to identify the scattered beam and target particles. The DSSSD contained four separate quadrants. Each consisted of 16 annular strips and 24 pairwise coupled radial strips. The results presented in this work were obtained by implementing the following conditions on the data: (1) building particle events based on a single front and back strip hit in the DSSSD; (2) selection of events in which one particle (1p) or two particles (2p) were detected; (3) applying a 2π time cut and requiring that the two particles were detected in opposite quadrants of the DSSSD; (4) identification of Sn and Ni particles through implementation of a kinematical cut; (5) reconstruction of the missing particles in the 1p data set; and (6) imposing a time coincidence condition between the particles and γ rays and subsequent event-by-event Doppler correction of the resulting γ-ray spectrum. The last condition is required as γ rays from the scattered beam and projectile particles are emitted in flight.

The Doppler-corrected spectra gated on Ni and Sn particles collected in the DSSSD are shown in panels (a) and (b) of Fig. 2, respectively. The data correspond to Ni particles collected in the angular range of 28.9–49.3°. A single peak, resulting from the 2⁺ → 0⁺ transition in 58Ni at 1454 keV, is seen in the Ni spectrum [21]. This state has an adopted reduced transition probability of B( E2; 0⁺ → 2⁺) = 0.0704(15) e²b² [21], which is used in the analysis described below. In the Sn spectrum, several previously known transitions are observed [22]. The extracted peak areas and relative intensities are given in Table I and the adopted level scheme for 109Sn is presented in Fig. 3 [22]. The peak around 664 keV is a combination of both the 3/2⁻⁰⁻→5/2⁺⁻⁻ and 5/2⁻⁻→7/2⁻⁺ transitions. Furthermore, it was not possible to resolve the 678-keV 5/2⁻⁻→5/2⁻⁺ transition. In the analysis, a fit consisting of two Gaussian peaks

Table I. Measured peak areas and relative intensities from the Coulomb excitation experiment. Energies given in Refs. [16,21,22].

<table>
<thead>
<tr>
<th>Transition</th>
<th>Eγ   (keV)</th>
<th>Area</th>
<th>Rel. int.</th>
</tr>
</thead>
<tbody>
<tr>
<td>Sn 3/2⁺⁻⁻→5/2⁺⁻⁻</td>
<td>665/664</td>
<td>462(40)</td>
<td>0.82(8)</td>
</tr>
<tr>
<td>/ 5/2⁻⁻→7/2⁻⁺</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Sn 5/2⁻⁻⁻→5/2⁺⁻⁻</td>
<td>678</td>
<td>60(31)</td>
<td>0.11(6)</td>
</tr>
<tr>
<td>Sn 3/2⁺⁻⁻→5/2⁺⁻⁻</td>
<td>925</td>
<td>277(32)</td>
<td>0.60(7)</td>
</tr>
<tr>
<td>Sn 1/2⁻⁻⁻→5/2⁻⁺</td>
<td>991</td>
<td>99(26)</td>
<td>0.22(6)</td>
</tr>
<tr>
<td>Sn 7/2⁺⁻⁻→5/2⁺⁻⁻</td>
<td>1078/1064</td>
<td>390(37)⁻</td>
<td>0.93(10)⁻</td>
</tr>
<tr>
<td>/ 7/2⁻⁺⁻→7/2⁻⁺⁻</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>Sn 9/2⁻⁻⁻→5/2⁺⁻⁻</td>
<td>1240</td>
<td>278(27)</td>
<td>0.72(8)</td>
</tr>
<tr>
<td>Ni 2⁺⁻⁻→0⁺⁻</td>
<td>1454</td>
<td>550(30)⁻</td>
<td>1.00(9)⁻</td>
</tr>
</tbody>
</table>

May also contain a contribution from the 3/2⁻⁻⁻→5/2⁺⁻⁻ transition at 1062 keV.

Calculated from the Ni peak area in Fig. 2(a) and the beam purity given in the text.
The analysis.

The labels indicate the transitions used in and (b) Sn particles. The spectra have been corrected for random peaks was used to extract areas for the 664-keV doublet and the 678-keV transition. The peak at 1078 keV may be a triplet related to the 1078-keV 7/2_1 → 5/2_1, 1064-keV 7/2_2 → 7/2_1, and 1062-keV 3/2_1 → 5/2_1 transitions. The weight of the peak is close to 1078 keV, suggesting that it is predominantly related to the 7/2_1 → 5/2_1 transition. In Ref. [15], the 7/2_1 → 7/2_1 transition was measured to be 22% of the intensity of the transition to the ground state. The peak area in Table I is therefore given as doublet containing these two contributions. The known branching ratios for the 678- and 1078-keV states were included in the analysis discussed below. The 991-keV level was not directly fed in previous β-decay experiments nor does it feed the 14-keV 7/2_1 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements. This procedure was carried out until the χ² + 1 limit was determined.

The shell-model value in the analysis [13]. During the fit, fifteen E2 and M1 matrix elements important for the reproduction of fifteen data points, including the measured yields and nine previously known branching ratios [16,22], were varied. The remaining matrix elements were fixed in the analysis. The determined transition probabilities are given in Table II. The errors were estimated by fixing each individual transition probability to points around the minimum. At each point, the fit was repeated while the Ni matrix elements were free to vary between their 1σ limits along with the remaining unfixed Sn matrix elements. This procedure was carried out until the χ² + 1 limit was determined. Each E2 transition was also fit individually with all other matrix elements set to zero. The effects are included in the total error,
TABLE II. Experimentally measured $B(E2)$ values in $^{109}$Sn compared to $B(E2)$ values from shell-model (SM) calculations and predictions based on a simple core-excitation (CE) model with a $^{108}$Sn core, $B(E2)=0.222(19)\ e^2b^2$ [7]. The SM values were used as starting values in the GOSIA2 analysis. For the $3/2^+$ states, both starting matrix elements tested in the analysis are given.

<table>
<thead>
<tr>
<th>$J_f$</th>
<th>$B(E2; 5/2^+ \rightarrow J_f)\ e^2b^2$</th>
<th>$\epsilon b^2$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$3/2^+$</td>
<td>0.029^{+0.014}_{-0.015}</td>
<td>0.007 or 0.001</td>
</tr>
<tr>
<td>$7/2^+$</td>
<td>0.060^{+0.020}_{-0.022}</td>
<td>0.047</td>
</tr>
<tr>
<td>$9/2^+$</td>
<td>0.085^{+0.032}_{-0.034}</td>
<td>0.032</td>
</tr>
<tr>
<td>$3/2^+$</td>
<td>0.013^{+0.009}_{-0.011}</td>
<td>0.001 or 0.007</td>
</tr>
<tr>
<td>$5/2^+$</td>
<td>$-0.012^a$</td>
<td>0.004</td>
</tr>
<tr>
<td>$1/2^+$</td>
<td>$-0.028^a$</td>
<td>0.008</td>
</tr>
</tbody>
</table>

$^a$Measured $\chi^2 + 1$ upper limits.

The $B(E2)$ values of three of the higher lying states are well described within the framework of the core-excitation model. The data suggest that these states belong to a core-excitation multiplet. The missing members are then the $1/2^+$ and $5/2^+$ levels. Two low-lying $1/2^+$ states have been observed in $^{109}$Sn. The lowest, at 545 keV, has been proposed to be the single-neutron state. The nonobservation of this level in the current measurement supports this assignment and suggests that the $1/2^+_2$ level belongs to the core-excitation multiplet. The $5/2^+_2$ state may be the missing $5/2^+$ member; however, its transition probability is largely overestimated by the core-excitation model. The shell-model gives a better description of the $E2$ properties of this state. Another possibility is that the $5/2^+$ multiplet member lies at a higher energy and was not directly excited in the experiment.

The $3/2^+_2$ state has been previously suggested to be the single-neutron state $^{[15,16]}$. The state is strongly populated in the $\beta$ decay of $^{109}$Sb, similar to the heavier isotopes $^{111-117}$Sn $^{[24]}$. The transition probabilities in Table II and shell-model calculations presented in this work, however, suggest that the $3/2^+_2$ state is a collective state. It may be that the strength of the collective excitation is distributed among both $3/2^+$ states. Interestingly, a low-lying $5/2^+$ state in $^{117}$Sn, suggested to be a single-neutron state from nucleon-transfer reactions, was measured in Coulomb excitation to have a transition probability of $\sim 6$ W.u. $^{[1]}$. The result was interpreted using the pairing-plus-quadrupole model and found to contain a large fraction of both the single-neutron and core-excited model states. A further consideration in $^{109}$Sn arises from the small energy difference between the $d_{5/2}$ and $g_{7/2}$ neutron orbits. Consequently, the low-lying excited states may also contain a fraction of the $g_{7/2}$ core-coupled state.

Recently, we reported on the measurements of the reduced transition probability for the lowest $3/2^+$ state in $^{107}$Sn $^{[3]}$. This state lies 40 keV in energy above the $3/2^+_1$ state in $^{109}$Sn, which suggests that they may have similar character. The measured $B(E2)$ value of $0.045^{+0.023}_{-0.017}\ e^2b^2$ for the lowest $3/2^+$ state in $^{107}$Sn is, however, $\sim 3$ times larger than the value for the $3/2^+_1$ state given in Table II. In $^{109}$Sn, the proposed single-neutron state is at 1280 keV $^{[16]}$. A possible explanation for the difference in the transition probabilities of the two low-lying $3/2^+$ states in these two nuclei is that the single-neutron state has moved up in energy in $^{107}$Sn and is less mixed with the collective state. The collective strength is instead concentrated on the one $3/2^+$ state in $^{107}$Sn as opposed to being distributed among the two $3/2^+$ states as in $^{109}$Sn.

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In conclusion, the measured transition probabilities can help to constrain the single-neutron energies relative to $^{100}$Sn. As an example, consider the transition probability of the $3/2^+_2$ state. As was discussed in Ref. $^{[3]}$, the inversion of the $g_{7/2}$ and $d_{5/2}$ orbits generally leads to a decrease in the calculated collectivity of the state. This effect can be compensated for by adjusting the energies of the three higher lying single-neutron orbits. The trend is that lowering the energy of $d_{3/2}$ state decreases the collectivity, while the opposite is true for the $s_{1/2}$ and $h_{11/2}$ orbits. However, the energy of the $h_{11/2}$ state is also constrained experimentally by the energies of the first
11/2− states in 107,109Sn. In the case discussed here, the measurements indicate that neither of the observed 3/2+ states is a d3/2 single-neutron state. The population of the 3/2+ states would not be expected if it was a pure single-neutron state, particularly since it lies ~300 keV higher than the 3/2− state. This observation effects the possibility of using the 3/2+ states in 109Sn to constrain the energy of the d3/2 orbit. On the other hand, the fact that we did not observe excitation of the 545-keV state strengthens the 1/2+ assignment of this state. It should also be mentioned that the shell-model calculations presented in this work do not include excitations across the N = Z = 50 closed proton and neutron shells. These missing excitations may increase the collectivity of the shell-model calculations; see Refs. [4–7].

We have presented the first Coulomb excitation measurement of the radioactive isotope 109Sn. A set of reduced transition probabilities for the observed excited states were deduced using the semi-classical Coulomb excitation code GOSIA2. The results indicate that a significant quadrupole strength exists in the low-lying excited states and thus highlight the importance of both single-neutron and collective excitations near the doubly magic nucleus 100Sn.

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