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Communication: THz absorption spectrum of the CO$_2$–H$_2$O complex: Observation and assignment of intermolecular van der Waals vibrations

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Terahertz absorption spectra have been recorded for the weakly bound CO$_2$–H$_2$O complex embedded in cryogenic neon matrices at 2.8 K. The three high-frequency van der Waals vibrational transitions associated with out-of-plane wagging, in-plane rocking, and torsional motion of the isotopic H$_2$O subunit have been assigned and provide crucial observables for benchmark theoretical descriptions of this system’s flat intermolecular potential energy surface. A (semi)-empirical value for the zero-point energy of 273 ± 15 cm$^{-1}$ from the class of intermolecular van der Waals vibrations is proposed and the combination with high-level quantum chemical calculations provides a value of 726 ± 15 cm$^{-1}$ for the dissociation energy $D_0$. © 2014 AIP Publishing LLC, [http://dx.doi.org/10.1063/1.4867901]

The intermolecular interaction between CO$_2$ and H$_2$O plays a major role for a variety of phenomena in physics, chemistry, and biology including the radiative transfer through the Earth’s and other planetary atmospheres, the early stages in carbonic acid formation and the transportation of dissolved CO$_2$ in the tissues of biological organisms. The prototypical binary mixed van der Waals complex of CO$_2$ and H$_2$O has been intensively investigated by both theory and experiment. Numerous quantum-chemical studies of the systems’ intermolecular potential energy surface (IPES) have been reported. High level of electron correlation, extensive basis sets, and inclusion of basis set superposition errors prove mandatory as the IPES is extremely flat near the global potential energy minimum. A recent comprehensive work reports a complete 5D $ab$ initio IPES composed of 23 000 high level single-point energies in configuration space and settles that the global potential energy minimum has a planar and T-shaped geometry of C$_{2v}$ symmetry with the oxygen atom of H$_2$O bound to the C atom and the H atoms pointing away from the CO$_2$ molecule (Fig. 1). This global potential energy minimum geometry has independently been confirmed experimentally by a variety of spectroscopic studies whereas other works have suggested the existence of other higher energy configurations. The most important observable is the dissociation energy $D_0$ which requires reliable band origins for the system’s complete set of fundamental vibrational transitions and in particular for the class of five intermolecular van der Waals vibrational transitions addressed in the present work.

Neon (L’Air Liquide, 99.5%) doped with CO$_2$ (Matheson, 99.9%) and degassed samples of H$_2$O, H$_2^{18}$O (Sigma Aldrich, 99.0% $^{18}$O) and D$_2$O (Sigma Aldrich, 99.5% D) with mixing ratios of $\approx$0.5% to 5%$_{eq}$ were deposited with a flow rate of 0.02 mol/h at 3.6 K on a gold-plated oxygen-free high thermal conductivity (OFHC) copper mirror inside an immersion helium cryostat (IHC-3) modified for matrix isolation spectroscopy at MAX-lab. A 3-mm deep cavity with a $\mu$m flat bottom has been drilled into the mirror centre to allow deposition of several millimeter thick matrices doped with weak absorbers. The mirror temperature was monitored by a Lake Shore silicon diode and maintained stable at 2.8 ± 0.1 K before and after deposition employing resistive heaters and feedback electronics. The sample mount was equipped with interchangeable CsI and poly-methylpentene (TPX) windows and combined IR and THz single-beam spectra were collected by a Bruker IFS 120 FTIR spectrometer employing a globar lamp as radiation source. A HgCdTe detector combined with a Ge/KBr beam splitter and a Si-bolometer operating at 4.2 K combined with a 6 $\mu$m multilayer Mylar beam splitter were employed for the IR and THz spectral regions, respectively. Spectral resolutions of 0.1–1 cm$^{-1}$ were selected depending on the observed bandwidths.

Figure 2 shows a series of THz absorption spectra collected for several millimeter thick cryogenic neon matrices doped with CO$_2$ (CO$_2$:H$_2$O:Ne=1:0.800), H$_2$O (0:1:4000), and mixtures of CO$_2$ and H$_2$O (3:1:4000), respectively. A

FIG. 1. The five intermolecular coordinates ($\alpha, \beta, \gamma, \theta$) specifying the configuration of the CO$_2$–H$_2$O van der Waals complex.
dominant spectral feature observed for the matrices doped with regular H$_2$O is observed at 79.5 cm$^{-1}$. This band has previously been assigned to a rotational-translation-coupling (RTC) transition of H$_2$O monomer in neon$^{21}$ and helps to monitor the H$_2$O monomer concentration. At higher H$_2$O concentrations, signs of hydrogen-bonded dimers of water started to show up. The strongest intermolecular hydrogen bond vibrations of (H$_2$O)$_2$ observed in the THz region, the acceptor torsional and acceptor wagging modes, have previously been observed and assigned in neon matrices at 116.0 cm$^{-1}$ and 122.2 cm$^{-1}$$^{20,22-24}$. THz spectra recorded for neon matrices doped solely with CO$_2$ showed no signs of CO$_2$-containing cluster entities in agreement with high-level theoretical studies of the weakly bound (CO$_2$)$_2$ system.$^{25}$ As shown in Fig. 2 the simultaneous use of CO$_2$ and H$_2$O as dopants allowed the identification of two new distinct bands located at 101.6 and 166.6 cm$^{-1}$; the latter band being a factor of 3–4 less intense. The intensity of these bands increased similarly with the CO$_2$ and H$_2$O concentration suggesting the assignments to a mixed (CO$_2$)$_n$–(H$_2$O)$_m$ complex. A stoichiometric 1:1 relationship could be confirmed at the low sample concentrations and by the complementary IR spectral series. THz spectra collected for neon matrices doped with CO$_2$ and isotopically enriched samples of H$_2^{18}$O and D$_2$O/HDO are useful to validate the proposed assignments further. An illustrative spectrum obtained for CO$_2$/H$_2^{18}$O doped neon matrices is shown in Fig. 2. The small isotopic shifts upon $^{18}$O-substitution have been reported previously both for the monomeric RTC transition$^{21}$ and the strongest dimeric hydrogen bond vibrations of water.$^{20}$ The observed spectral shifts caused by the $^{18}$O-substitution are also very small for both the proposed CO$_2$–H$_2$O bands. The band at 101.6 cm$^{-1}$ seems to be rather unaffected by the $^{18}$O-substitution. The upper band has a small but significant reproducible isotopic red-shift of 1.9 cm$^{-1}$ as indicated by arrows in Fig. 2. The effect of deuteration has a much larger impact on the recorded THz spectra since both the monomeric and dimeric water spectral features have significant isotope shifts not to mention that mixed isotopic entities start to build up. Both the D$_2$O and HDO as well as the mixed deuterated dimeric entities of water have been extensively studied previously and the complete list of spectral assignments is given in Table I. In the further spectral analysis of the isotopically enriched spectral signatures we consult quantum chemical predictions.

As described and cited in the introduction, a variety of high-level quantum chemical calculations have been performed to describe the accurate IPES of the CO$_2$–H$_2$O complex in great detail. In Table II, we have listed harmonic MP2/aug-cc-pVQZ predictions for the global potential energy
minimum with \( C_{2v} \) symmetry. These predictions are in accordance with the study by Makarewicz\(^{14} \) and specify the class of five intermolecular van der Waals vibrational transitions. The out-of-plane H\(_2\)O wagging mode \( \nu_8 \) (\( B_1 \)) described by the intermolecular coordinate \( \beta \) (Fig. 1) has a high intensity whereas the in-plane rocking of H\(_2\)O \( \nu_{11} \) (\( B_2 \)) described by \( \alpha \) and the CO\(_2\) librational mode of \( \nu_1 \) (\( B_2 \)) described by \( \theta \) both have medium intensity. The H\(_2\)O torsional mode \( \nu_6 \) (\( A_2 \)) described by \( \gamma \) is strictly IR-forbidden for the regular CO\(_2\)–H\(_2\)O complex and the intermolecular \( O\cdots C \) stretching mode \( \nu_3 \) (\( A_1 \)) described by \( r \) has a very low intensity. A convincing correlation immediately appears between the harmonic predictions and the observation of the two distinct bands at 101.6 and 166.6 cm\(^{-1} \) for THz spectra of neon matrices doped simultaneously with CO\(_2\) and H\(_2\)O. This excellent agreement between experiment and theory suggests straightforward assignments for the out-of-plane H\(_2\)O wagging mode \( \nu_8 \) at 101.6 cm\(^{-1} \) and the in-plane H\(_2\)O rocking mode \( \nu_{11} \) at 166.6 cm\(^{-1} \) for the regular CO\(_2\)–H\(_2\)O complex. This agreement could partly be due to some kind of fortunate cancelation of smaller spectral shifts in opposite directions originating from anharmonicity effects and minor matrix perturbations which shall be discussed below. Small isotopic red-shifts are expected upon \(^{18}\)O-substitution on the water subunit since both these intermolecular vibrational modes involve hindered rotational motion of H\(_2\)O. The harmonic calculations for the enriched CO\(_2\)–H\(_2\)\(^{18}\)O system predict a very small red-shift of 0.9 cm\(^{-1} \) for the H\(_2\)O wagging mode \( \nu_8 \) and a more pronounced red-shift of 2.2 cm\(^{-1} \) for the H\(_2\)O rocking mode \( \nu_{11} \) (Table III). These red-shift predictions are convincingly close to the observations within the experimental reproducibility. The predicted relative harmonic band intensity agrees qualitatively with the experimental findings. The rotational motion of D\(_2\)O is much slower relative to H\(_2\)O and rather large isotopic spectral red-shifts for both the D\(_2\)O wagging and rocking modes are expected. The harmonic calculations for CO\(_2\)–D\(_2\)O accordingly predict red-shifts of 24.3 and 19.1 cm\(^{-1} \) with factors of 2 and 3 smaller harmonic band intensities, respectively, for these intermolecular van der Waals modes (Table III). The band origin for the strongest D\(_2\)O wagging mode is expected at 77.4 cm\(^{-1} \) and thereby unavoidably overlapped with the strong RTC transitions for H\(_2\)O, HDO, and D\(_2\)O (Table I). The band origin for the D\(_2\)O rocking mode with medium intensity is predicted at 148.5 cm\(^{-1} \) but the threefold loss of band intensity (8.5 km/mol) and the experimental challenge with isotopic H/D exchange and formation of HDO in the matrix inlet system was expected to blur this spectral signature. Nevertheless, a series of THz spectra recorded for neon matrices doped with CO\(_2\)/HDO/D\(_2\)O mixtures reproduces a medium strong band at 145.1 cm\(^{-1} \) which we assign to the D\(_2\)O rocking transition of the CO\(_2\)–D\(_2\)O complex as indicated in Fig. 2. The rotational motion of HDO is slower relative to H\(_2\)O but faster relative to D\(_2\)O and smaller predicted isotopic spectral red-shifts for the H\(_2\)O wagging and rocking modes of 18.1 and 11.6 cm\(^{-1} \), respectively, for CO\(_2\)–HDO relative to CO\(_2\)–H\(_2\)O are both expected. The harmonic band intensities are both predicted to be a factor of 2 smaller for this isotopic variant of the complex. The band origin for the stronger HDO wagging mode predicted at 83.6 cm\(^{-1} \) cannot be observed unambiguously due to the overlapping RTC transitions of H\(_2\)O, HDO, and D\(_2\)O. The red-shifted HDO rocking mode predicted at 156.0 cm\(^{-1} \) with 14.5 km/mol intensity is easily observed and assigned at 152.1 cm\(^{-1} \) for the CO\(_2\)/HDO/D\(_2\)O doped neon matrices.

### Table I. Assignment of transitions (units of cm\(^{-1} \)) for the recorded THz absorption spectra of CO\(_2\)/H\(_2\)\(^{16}\)O, CO\(_2\)/H\(_2\)\(^{18}\)O, CO\(_2\)/HDO, and CO\(_2\)/D\(_2\)O embedded in neon matrices at 2.8 K.

<table>
<thead>
<tr>
<th>( \text{CO}_2/\text{H}_2^{16}\text{O} )</th>
<th>( \text{CO}_2/\text{H}_2^{18}\text{O} )</th>
<th>( \text{CO}_2/\text{HDO} )</th>
<th>( \text{CO}_2/\text{D}_2\text{O} )</th>
<th>Assignment</th>
</tr>
</thead>
<tbody>
<tr>
<td>79.5(^{a})</td>
<td>77.8(^{a})</td>
<td>73.1(^{a})</td>
<td>H(_2)O monomer, RTC transition(^{14})</td>
<td></td>
</tr>
<tr>
<td>101.6(^{b})</td>
<td>101.4(^{b})</td>
<td></td>
<td>CO(_2)-H(_2)O complex, H(_2)O wagging</td>
<td></td>
</tr>
<tr>
<td>116.0(^{c})</td>
<td>116.0(^{c})</td>
<td>(H(_2)O)(_2) complex, acceptor torsion</td>
<td></td>
<td></td>
</tr>
<tr>
<td>122.2(^{c})</td>
<td>122.0(^{c})</td>
<td>106.0(^{c})</td>
<td>(H(_2)O)(_2) complex, acceptor wagging</td>
<td></td>
</tr>
<tr>
<td>150.6(^{b})</td>
<td>149.7(^{b})</td>
<td>135.0(^{b})</td>
<td>CO(_2)-H(_2)O complex, H(_2)O torsion</td>
<td></td>
</tr>
<tr>
<td>166.6(^{b})</td>
<td>164.7(^{b})</td>
<td>152.1(^{b})</td>
<td>(H(_2)O)(_2) complex, acceptor twist</td>
<td></td>
</tr>
</tbody>
</table>

\(^{a}\)Rotation-translation-coupling transition (see Ref. 21).

\(^{b}\)Present work.

\(^{c}\)Reference 20.

### Table II. MP2/aug-cc-pVQZ predictions in the double harmonic approximation of vibrational band origins (units of cm\(^{-1} \)) and corresponding infrared band strengths (units of km/mol, in parenthesis) for the weakly bound CO\(_2\)-H\(_2\)O complex of \( C_{2v} \) symmetry.

<table>
<thead>
<tr>
<th>Mode ( \nu_1 ) (( A_1 ))</th>
<th>Description</th>
<th>Origin (Int.)</th>
<th>Mode ( \nu_7 ) (( B_1 ))</th>
<th>Description</th>
<th>Origin (Int.)</th>
</tr>
</thead>
<tbody>
<tr>
<td>( v_1 ) (( A_1 ))</td>
<td>Sym. H(_2)O stretch</td>
<td>3836.3 (10)</td>
<td>( v_7 ) (( B_1 ))</td>
<td>Out-of-plane CO(_2) bend</td>
<td>667.2 (22)</td>
</tr>
<tr>
<td>( v_2 ) (( A_1 ))</td>
<td>In-plane H(_2)O bend</td>
<td>1630.2 (76)</td>
<td>( v_8 ) (( B_1 ))</td>
<td>Internol. H(_2)O wagging</td>
<td>101.7 (228)</td>
</tr>
<tr>
<td>( v_3 ) (( A_1 ))</td>
<td>Sym. CO(_2) stretch</td>
<td>1334.7 (0.2)</td>
<td>( v_9 ) (( B_2 ))</td>
<td>Antisym. H(_2)O stretch</td>
<td>3962.3 (87)</td>
</tr>
<tr>
<td>( v_4 ) (( A_1 ))</td>
<td>In-plane CO(_2) bend</td>
<td>653.2 (39)</td>
<td>( v_{10} ) (( B_2 ))</td>
<td>Antisym. CO(_2) stretch</td>
<td>2415.3 (364)</td>
</tr>
<tr>
<td>( v_5 ) (( A_1 ))</td>
<td>Internol. O\cdots C stretch</td>
<td>111.2 (0.4)</td>
<td>( v_{11} ) (( B_2 ))</td>
<td>Internol. H(_2)O rocking</td>
<td>167.6 (25)</td>
</tr>
<tr>
<td>( v_6 ) (( A_2 ))</td>
<td>Internol. H(_2)O torsion</td>
<td>151.2 (0.0)</td>
<td>( v_{12} ) (( B_2 ))</td>
<td>Internol. CO(_2) libration</td>
<td>16.8 (29)</td>
</tr>
</tbody>
</table>
This assignment is indicated in Fig. 2 for an experiment with ν

The contributions from the H2O wagging mode ν6 observed at 101.6 cm⁻¹ and the H2O rocking mode ν11 observed at 166.6 cm⁻¹ are clear. In order to establish the contribution from the IR-forbidden H2O torsional mode ν6 we consider the effect of H/D substitution on this kind of motion. The torsional mode can be described as a hindered rotation of H2O around the C2 symmetry axis which also has the role of the principal b-axis for an isolated H2O molecule. The torsional mode can as such be described as an almost genuine b-type rotation of H2O although this motion is hindered in the complex. In the harmonic approximation the ratio of the two harmonic torsional band origins is thus proportional to the square root of the ratio of the B rotational constants for the two isotopic water molecules in question.

In this way, the use of experimental B-values predicts an isotopic harmonic red-shift of 21% for the band origin upon a single H/D-substitution. The model predicts a blue-shift of the proposed HDO torsional band origin at 128.4 cm⁻¹ to 162.1 cm⁻¹ for the corresponding IR-forbidden mode of the regular CO2–H2O system where our MP2 predictions provide a band origin close of 151.2 cm⁻¹.

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